Noncommutative polynomial optimization and Bell inequalities

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1 Motivation: Linear Bell inequalities

2 Noncommutative polynomial optimization and the Navascués-Pironio-Acín hierarchy

3 Nonlinear Bell inequalities and state polynomial optimization

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Classical correlations



 $\mathsf{Alice} \gets \mathsf{Source} \to \mathsf{Bob}$

• Correlations = conditional joint probabilities:

P(ab | st) = P(Alice, Bob answer a, b | Alice, Bob are asked s, t)

• Deterministic strategies:

 $P(a \mid s), P(b \mid t) \in \{0, 1\} \Rightarrow P(ab \mid st) = P(a \mid s)P(b \mid t) \in \{0, 1\}$

Classical correlations: convex comb. of deterministic correlations

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- Linear Bell inequalities: linear inequalities in terms of correlations (e.g., P(ab | st)) & marginals (e.g., P(a | s), P(b | t)), valid for all classical correlations
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• Clauser-Horne-Shimony-Holt (CHSH) inequality:

 $P(11 \mid 00) + P(11 \mid 01) + P(11 \mid 10) - P(11 \mid 11) \le 2$

• Alice & Bob share a bipartite quantum state ψ and they answer a, b by performing quantum measurements on their part of ψ :

 $P(ab \mid st) = \psi^* A^a_s B^b_t \psi$

for some commuting unitary operators $\{A_s^a\}, \{B_t^b\}$

• CHSH: $\psi^*(A_0^1B_0^1 + A_0^1B_1^1 + A_1^1B_0^1 - A_1^1B_1^1)\psi \le 2$ is violated on the entangled state $\psi = \frac{1}{\sqrt{2}}(e_1 \otimes e_1 + e_2 \otimes e_2)$ (attaining the value $2\sqrt{2}$)

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- Question: Given a linear Bell inequality, how could we determine the maximal quantum violation?
- Reformulate as an optimization problem in operators:

$$\xi = \sup \{ \langle f(A, B)\psi, \psi \rangle : A_i^2 = 1, B_j^2 = 1, A_i B_j = B_j A_i, \psi \in \mathcal{H}, \|\psi\| = 1 \}$$

= inf λ s.t. $\lambda I - f(A, B) \succeq 0$, for all $A_i^2 = 1, B_j^2 = 1, A_i B_j = B_j A_i$

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Noncommutative polynomial optimization

• $\mathcal{B}(\mathcal{H})$: Banach algebra of bounded linear operators on \mathcal{H}

• Words/monomials: products of $x_1, \ldots, x_n \in \mathcal{B}(\mathcal{H})$

• Noncommutative (NC) polynomials: polynomials in $x_1, \ldots, x_n \in \mathcal{B}(\mathcal{H})$

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$$\xi := \begin{cases} \inf_{\mathbf{x} \in \mathcal{B}(\mathcal{H})^n} & \lambda_{\min}(f(\mathbf{x})) \ (= \langle \psi | f(\mathbf{x}) | \psi \rangle \text{ where } \psi \in \mathcal{H}, \|\psi\| = 1) \\ \text{s.t.} & g_i(\mathbf{x}) \succeq 0, \quad i = 1, \dots, m \end{cases}$$

The Navascués-Pironio-Acín hierarchy of semidefinite relaxations

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Reformulation with nonnegative NC polynomials

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$$\mathcal{D} := \left\{ \mathbf{x} \in \mathcal{B}(\mathcal{H})^n \, \Big| \, g_i(\mathbf{x}) \succeq 0, \, i = 1, \dots, m \right\}$$

$$\xi = \sup_{\lambda} \left\{ \lambda : f(\mathbf{x}) - \lambda I \succeq \mathbf{0}, \, \forall \mathbf{x} \in \mathcal{D} \right\}$$

 \bullet Problem: characterize the set of NC polynomials that are positive semidefinite over $\mathcal D$

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• Sum of Hermitian squares (SOHS): $f = f_1^* f_1 + \cdots + f_t^* f_t$

• $[\mathbf{x}]_d \coloneqq [1, x_1, \dots, x_n, x_1^d, \dots, x_n^d]$

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Theorem (Helton and McCullough, 2004)

Assume that $\{g_i\}_{i=1}^m$ satisfies Archimedean's condition. If f is positive on

 \mathcal{D} , then

$$f = \sum_{g \in \{1\} \cup \{g_i\}_{i=1}^m} s_g^* g s_g, \text{ for some } \{s_g\}_g \subseteq \mathbb{R}[\mathsf{x}].$$
• *r*-th order SOHS relaxation:

$$\theta_r^* \coloneqq \begin{cases} \sup_{\lambda, s_g} & \lambda \\ \text{s.t.} & f - \lambda = \sum_{g \in \{1\} \cup \{g_i\}_{i=1}^m} s_g^* g s_g & \rightsquigarrow & \text{SDP} \\ & \deg(s_g^* g s_g) \le 2r \end{cases}$$

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 $\mathfrak{P} \theta_r^* \nearrow \xi$ as $r \to \infty$ under Archimedean's condition

Consider a linear functional L : ℝ[x] → ℝ such that L(w*) = L(w)
 r-th order moment matrix M_r(L):

$$[\mathbf{M}_r(L)]_{uv} \coloneqq L(u^*v), \quad \forall |u|, |v| \le r$$

• *r*-th order localizing matrix $M_r(gL)$ associated to $g = \sum_w g_w w$:

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The Navascués-Pironio-Acín hierarchy



Noncommutative polynomial optimization

Detecting global optimality and extracting an optimal solution

• Global optimality is certified (i.e., $\theta_r = \xi$) if

$\operatorname{rank} M_r(L) = \operatorname{rank} M_{r-d_D}(L)$ (flatness condition)

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• The size of SDPs arising from the hierarchy rapidly grows

• Exploiting structures:

- Using a customized monomial basis
- Equality constraints

Symmetry

- Correlative sparsity
- Term sparsity

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Ground state energy of quantum many-body systems

• The Heisenberg chain is defined by the Hamiltonian:

$$H = \sum_{i=1}^{N} \sum_{a \in \{x, y, z\}} \sigma_i^a \sigma_{i+1}^a$$

• Computing ground state energy of the Heisenberg chain can be

formulated as a noncommutative polynomial optimization problem:

$$\begin{split} \min_{\{|\psi\rangle,\sigma_i^a\}} & \langle \psi|H|\psi\rangle \\ \text{s.t.} & (\sigma_i^a)^2 = 1, \quad i = 1, \dots, N, a \in \{x, y, z\} \\ & \sigma_i^x \sigma_i^y = \mathbf{i}\sigma_i^z, \sigma_i^y \sigma_i^z = \mathbf{i}\sigma_i^x, \sigma_i^z \sigma_i^x = \mathbf{i}\sigma_i^y, \quad i = 1, \dots, N \\ & \sigma_i^a \sigma_j^b = \sigma_j^b \sigma_i^a, \quad 1 \le i \ne j \le N, a, b \in \{x, y, z\} \end{split}$$

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Structures of the Heisenberg model

sparsity

- e sign symmetry
- Itranslation symmetry
- opermutation symmetry
- In mirror symmetry

100 spins, accuracy 10^{-5} [Wang et al., PRX 2024]

From linearity to nonlinearity

Covariance Bell inequalities

• Binary random variables A and B

• Covariance on a classical model:

$$\operatorname{cov}(A,B) = \int AB d\mu - \int A d\mu \int B d\mu = E(AB) - E(A)E(B)$$

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• Covariance on a quantum model:

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$$cov_{3322} = cov(A_1, B_1) + cov(A_1, B_2) + cov(A_1, B_3) + cov(A_2, B_1) + cov(A_2, B_2) - cov(A_2, B_3) + cov(A_3, B_1) - cov(A_3, B_2)$$

> Concrete
$$\mu$$
 yields $\cos_{3322} = 4.5$

$$\mu = \frac{3}{8}(+++/+++) + \frac{3}{8}(--+/--+) + \frac{1}{4}(-+-/-+-)$$

► $(A_1A_2A_3/B_1B_2B_3)$: strategy where Alice and Bob deterministically output A_x and B_y for inputs x and y

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► $(A_1A_2A_3/B_1B_2B_3)$: strategy where Alice and Bob deterministically output A_x and B_y for inputs x and y

> Concrete ς yields $cov_{3322} = 5$

What is the maximal quantum violation?



 $\mathsf{Alice} \gets \mathsf{Source} \ 1 \to \mathsf{Bob} \gets \mathsf{Source} \ 2 \to \mathsf{Charlie}$

• Observers hold particles from different sources and therefore a priori share no correlations

• A party that holds multiple shares originating from different sources can perform entangled measurements to a posteriori distribute entanglement



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• Binary random variables A_i, B_j, C_k :

$$\frac{1}{3}\sum_{i\in\{1,2,3\}} \left(E(B_iC_i) - E(A_iB_i) \right) - \sum_{\{i,j,k\}=\{1,2,3\}} E(A_iB_jC_k)$$

• Bilocality constraints $E(A_1A_2A_3C_1C_2C_3) = E(A_1A_2A_3)E(C_1C_2C_3) +$

similar factorization constraints & vanishing constraints:

 $E(A_i) = E(B_i) = E(C_i) = 0$ for $i \in \{1, 2, 3\}$

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Jie Wang

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State polynomial optimization

- States $\mathcal{S}(\mathcal{H})$: positive unital \star -linear functionals on $\mathcal{B}(\mathcal{H})$
- State polynomial: $\mathscr{S} \coloneqq \mathbb{R}[\varsigma(w) \colon w \in [\mathbf{x}] \setminus \{1\}]$ $\succ \varsigma(x_1^2)\varsigma(x_2x_1) + \varsigma(x_1)\varsigma(x_2x_1x_2), x_1, \dots, x_n \in \mathcal{B}(\mathcal{H}), \varsigma \in \mathcal{S}(\mathcal{H})$
- NC state polynomial: $\mathscr{S} := \mathscr{S} \otimes \mathbb{R}[\mathbf{x}]$, the free \mathscr{S} -algebra on $[\mathbf{x}]$ > $\varsigma(x_1^2)x_2x_1 + \varsigma(x_1)\varsigma(x_2x_1x_2), x_1, \dots, x_n \in \mathcal{B}(\mathcal{H}), \varsigma \in \mathcal{S}(\mathcal{H})$

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- NC state polynomial: 𝒴 := 𝒴 ⊗ ℝ[x], the free 𝒴-algebra on [x]

 ς(x₁²)x₂x₁ + ζ(x₁)ζ(x₂x₁x₂), x₁,...,x_n ∈ 𝔅(𝔑), ζ ∈ 𝔅(𝔑)

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•
$$f \in \mathscr{S}, g_i \in \mathscr{S}$$
:

$$\begin{cases} \inf_{\substack{(\varsigma, \mathbf{x}) \\ \text{s.t.} \quad g_i(\varsigma; \mathbf{x}) \succeq 0, \quad i = 1, \dots, m \end{cases}} \end{cases}$$

- Sum of Hermitian squares (SOHS): $f = \varsigma(f_1^* f_1) + \cdots + \varsigma(f_t^* f_t)$
- NC state monomial: $\varsigma(u_1) \cdots \varsigma(u_s) w$
- $\langle \mathbf{x} \rangle_d$: vector of NC state monomials up to degree d
- f of degree 2d is an SOHS \iff there exists a PSD matrix G s.t. $f = \varsigma \left(\langle \mathbf{x} \rangle_d^* \cdot G \cdot \langle \mathbf{x} \rangle_d \right)$

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Define the feasible set

$$\mathcal{D} \coloneqq \left\{ (\varsigma, \mathbf{x}) \in \mathcal{S}(\mathcal{H}) \times \mathcal{B}(\mathcal{H})^n \, \middle| \, g_i(\varsigma; \mathbf{x}) \succeq 0, \, i = 1, \dots, m \right\}$$

Theorem (Klep, Magron, Volčič, and Wang, 2024)

Assume that $\{g_i\}_{i=1}^m$ is algebraically bounded. If f is positive on \mathcal{D} , then

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Theorem (Klep, Magron, Volčič, and Wang, 2024)

Assume that $\{g_i\}_{i=1}^m$ is algebraically bounded. If f is positive on \mathcal{D} , then

$$f = \sum_{g \in \{1\} \cup \{g_i\}_{i=1}^m} \varsigma(s_g^* g s_g), \text{ for some } \{s_g\}_g \subseteq \mathscr{S}.$$

• *r*-th order SOHS relaxation:

$$\theta_r^* := \begin{cases} \sup_{\lambda, s_g} & \lambda \\ \text{s.t.} & f - \lambda = \sum_{g \in \{1\} \cup \{g_i\}_{i=1}^m} \varsigma(s_g^* g s_g) & \rightsquigarrow & \text{SDP} \\ & \deg(s_g^* g s_g) \le 2r \end{cases}$$

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 $\mathfrak{P} \theta_r^* \nearrow \xi$ as $r \to \infty$ under the algebraically bounded condition

• Consider a linear functional $L : \mathscr{S} \to \mathbb{R}$

• *r*-th order moment matrix $M_r(L)$:

$$[\mathbf{M}_r(L)]_{uv} \coloneqq L(\varsigma(u^*v)), \quad \forall |u|, |v| \le r$$

• *r*-th order localizing matrix $M_r(gL)$ associated to $g = \sum_w g_w w$:

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• $E(A_2 + B_1 + B_2 - A_1B_1 + A_2B_1 + A_1B_2 + A_2B_2) - E(A_1)E(B_1) - E(A_2)E(B_1) - E(A_2)E(B_2) - E(A_1)^2 - E(B_2)^2$

• For classical models: 3.375

$$\begin{aligned} \sup_{x_i, y_j} & \varsigma(x_2) + \varsigma(y_1) + \varsigma(y_2) - \varsigma(x_1y_1) + \varsigma(x_2y_1) + \varsigma(x_1y_2) + \varsigma(x_2y_2) \\ & -\varsigma(x_1)\varsigma(y_1) - \varsigma(x_2)\varsigma(y_1) - \varsigma(x_2)\varsigma(y_2) - \varsigma(x_1)^2 - \varsigma(y_2)^2 \\ \text{s.t.} & x_i^2 = 1, y_j^2 = 1, [x_i, y_j] = 0 \text{ for } i, j = 1, 2 \end{aligned}$$

• For quantum models: 3.5114 (r = 2)

- $E(A_2 + B_1 + B_2 A_1B_1 + A_2B_1 + A_1B_2 + A_2B_2) E(A_1)E(B_1) E(A_2)E(B_1) E(A_2)E(B_2) E(A_1)^2 E(B_2)^2$
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$$\begin{aligned} \sup & \frac{1}{3} \sum_{i \in \{1,2,3\}} \left(\varsigma(b_i c_i) - \varsigma(a_i b_i) \right) - \sum_{\{i,j,k\} = \{1,2,3\}} \varsigma(a_i b_j c_k) \\ \text{s.t.} & \varsigma(a_1 a_2 a_3 c_1 c_2 c_3) = \varsigma(a_1 a_2 a_3) \varsigma(c_1 c_2 c_3) \\ & \varsigma(a_i) = \varsigma(b_i) = \varsigma(c_i) = 0 \text{ for } i \in \{1,2,3\} \\ & \varsigma(a_i b_j) = \varsigma(b_i c_j) = 0 \text{ for } i \neq j \\ & \varsigma(a_i b_j c_k) = 0 \text{ for } |\{i,j,k\}| \le 2 \end{aligned}$$

• SDP with r = 3 yields 4.46 (few seconds)

• SDP with r = 4 yields 4.38 (few hours)

• SDP with r = 5 yields 4.37 (one week)

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- Quantum uncertainty relation [Morán and Huber, 2024]
- Quantum code [Munné, Nemec, and Huber, 2024]
- Quantum causal compatibility problem [Ligthart, 2024]
- Quantum nonlocality [Renou, Xu, and Ligthart, 2024]

• NCTSSOS: in Julia, based on JuMP

https://github.com/wangjie212/NCTSSOS

• Play with it if you have noncommutative/state polynomial optimization problems!

Noncommutative polynomial optimization is a powerful tool for handling nonlinear expressions in operators arising in quantum information State polynomial optimization is a powerful tool for handling nonlinear expressions in both expectations and operators arising in quantum

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Thank You!

https://wangjie212.github.io/jiewang